## AOM-HW3 with Code

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#### September 30, 2020

I display relevant portions of code throughout the document. At the very end of the document, I include the entire code in one piece - See Appendix A

### 1 Problem 1

Most of the explanation for problem 1 was written on a sheet of paper, which is attached to the submitted pdf document.

Here is the code for the COE2RV and RV2COE functions:

```
def COE2RV(a, e, i, RAAN, w, ta):
  1
               # Used https://web.archive.org/web/20160418175843/https://ccar.colorado.edu/
                  asen5070/handouts/cart2kep2002.pdf as a reference
 3
               r = (a * (1 - e * 2)) / (1 + e * np.cos(ta)) # Get position from orbit formula
  4
               h = np.sqrt(mu * a * (1 - e * 2)) \# Magnitude of specific angular momentum
 5
               X = r * (np.cos(RAAN) * np.cos(w + ta) - np.sin(RAAN) * np.sin(w + ta) * np.cos(i)
 6
               ) # Position X-component
Y = r * (np.sin(RAAN) * np.cos(w + ta) + np.cos(RAAN) * np.sin(w + ta) * np.cos(i
                 )) # Position Y-component
               Z = r * (np.sin(i) * np.sin(w + ta)) # Position Z-component
 8
               p = a * (1 - e * 2) \# Semilatus Rectum
 9
               X_{dot} = ((X * h * e) / (r * p)) * np.sin(ta) - ((h / r) * (np.cos(RAAN) * np.sin(w)))
                  + ta) + np.sin(RAAN) * np.cos(w + ta) * np.cos(i))) # Velocity X-component
               Y_{dot} = ((Y * h * e) / (r * p)) * np.sin(ta) - ((h / r) * (np.sin(RAAN) * np.sin(w range))) + (r * p)) * np.sin(ta) - ((h / r) * (np.sin(RAAN) * np.sin(w range))) + (r * p)) * np.sin(ta) - ((h / r) * (np.sin(RAAN) * np.sin(w range))) + (r * p)) * np.sin(ta) - ((h / r) * (np.sin(RAAN) * np.sin(w range))) + (r * p)) * np.sin(ta) - ((h / r) * (np.sin(RAAN) * np.sin(w range))) + (r * p)) * np.sin(ta) - ((h / r) * (np.sin(RAAN) * np.sin(w range))) + (r * p)) * np.sin(w range)) + (r * p)) + (r * p))
                  + ta) - np.cos(RAAN) * np.cos(w + ta) * np.cos(i))) # Velocity Y-component
               Z_{dot} = ((Z * h * e) / (r * p)) * np.sin(ta) + ((h / r) * (np.sin(i) * np.cos(w + ta)))
                  )) # Velocity Z-component
13
               r_vector_from_COE2RV = np.array([X, Y, Z]) # Put X,Y,Z into an array to create a
14
                 vector
               v\_vector\_from\_COE2RV = np.array([X\_dot, Y\_dot, Z\_dot]) # Put X\_dot, Y\_dot, Z\_dot]
                  into an array to create a vector
               #print("Position Vector (km)= " + str(r_vector_from_COE2RV))
               \#print("Velocity Vector (km/s) = " + str(v_vector_from_COE2RV))
18
19
               return(r_vector_from_COE2RV, v_vector_from_COE2RV)
20
21
      def RV2COE(r_vector, v_vector):
22
               # Used "Orbital Mechanics for Engineering Students" 3rd Edition by Howard D.
23
                  Curtis; Also used our class notes
24
```

```
r = np.linalg.norm(r_vector) \# Distance
25
       v = np.linalg.norm(v_vector) # MAgnitude of Velocity or Speed
26
       v_r = np.dot(v_vector, r_vector) / r # radial velocity
27
       h_vector = np.cross(r_vector, v_vector) # specific angular momentum vector
28
       h = np.linalg.norm(h_vector) # magnitude of specific angular momentum
29
       i = np.arccos(h_vector [2] / h) # inclination
n_vector = np.cross([0, 0, 1], h_vector) # vector pointing to ascending node
30
31
       n = np.linalg.norm(n_vector) # magnitude of n
32
       if n_vector[1] > 0:
33
           RAAN = np.arccos(n_vector[0] / n) \# Right Ascension of the Ascending node
34
35
       if n_vector [1] < 0:
           RAAN = 2 * np.pi - np.arccos(n_vector[0] / n)
36
       e_vector = (1 / mu) * ((((v ** 2) - (mu / r)) * (r_vector)) - ((r * v_r) * v_vector))
37
         \# eccentricity vector
       e = np.linalg.norm(e_vector) \# eccentricity
38
       if e_vector [2] > 0:
39
40
           w = np.arccos(np.dot(n_vector, e_vector) / (n * e)) # Argument of periapse
       if e_vector [2] < 0:
41
           w = 2 * np.pi - np.arccos(np.dot(n_vector, e_vector) / (n * e))
42
       if v_r > 0:
43
           ta = np.arccos((np.dot(e_vector, r_vector) / (e * r))) # True anomaly
44
       if v_r < 0:
45
           ta = 2 * pi - np.arccos((r_vector / r) * (e_vector / e))
46
47
       Energy = (v ** 2 / 2) - (mu / r)
       if e == 1:
48
           p = h * 2 / mu \# Semilatus Rectum
49
           return "Orbit is parabolic. Eccentricity = infinity"
50
51
       else:
           a = -mu / (2 * Energy) # Semi-major Axis
           p = a * (1 - e * 2)
53
       print("Semi-major Axis (km) = " + str(a))
55
       print("Eccentricity = " + str(e))
56
       print("Inclination (deg)= " + str(i * (180 / np.pi)))
57
       print("Right Ascension of the Ascending Node (deg) = " + str(RAAN * (180 / np.pi)))
58
       print("Argument of Periapse (deg)= " + str(w * (180 / np.pi)))
59
       print ("True Anomaly (deg) = " + str(ta * (180 / np.pi)))
60
61
       return [a, e, i, RAAN, w, ta]
62
63
64
mu = 398600 \# Gravitational Parameter of Earth km<sup>3</sup>/s<sup>2</sup>
66 r_vector = np.array([-6045.0, -3490.0, 2500.0])
   v_vector = np.array([-3.457, 6.618, 2.533])
67
68
69
   COE = RV2COE(r_vector, v_vector)
70
71 COE2RV(COE[0], COE[1], COE[2], COE[3], COE[4], COE[5])
```

The second to last line of code converts the position and velocity vectors to its classical orbital elements. The last line of code takes those elements and turns them back into the position and velocity vectors. I was returned the original position and velocity vectors back exactly.

### 2 Problem 2

To propagate the frozen orbit, from the previous assignment, I wrote my own Runge-Kutta 4th order (RK4) integrator that numerically integrates the unperturbed two-body equations of motion. We can break down the two-body equation of motion

$$\ddot{\mathbf{r}} = \frac{-\mu}{r^3} \mathbf{r} \tag{1}$$

into two differential equations with only one derivative in each:

$$\frac{d\mathbf{r_i}}{dt} = \mathbf{v_i} \tag{2}$$

$$\frac{d\mathbf{v_i}}{dt} = \frac{-\mu}{r^3} \mathbf{r_i} \tag{3}$$

where  $\mathbf{r}_i$  is the position vector of body i (i = 1, 2),  $\mathbf{v}_i$  is the velocity vector of body i, and r is the norm of the vectors ( $\mathbf{r}_2 - \mathbf{r}_1$ ). Because our satellite is orbiting Earth in a geocentric reference frame, we can set the position and velocity vectors of Earth (the first body) equal to 0.

I created a function named "force" that calculates eq. 2 and eq. 3 for a given  ${\bf r}$  vector

1	# Calculate State Vectors r & v using the COE2RV I created above
2	$r_sat$ , $v_sat = COE2RV(a, e, i, raan, w, ta) # Position state vector (km) and$
	Velocity state vector (km/s) of satellite in the geocentric equitorial frame (km)
3	
4	t = 5 * T #Total time for propagation
5	dt = 1 # time - step
6	$t_{array} = np.linspace(0, t, t / dt + 1) #make a time array$
7	$X = np.array([r\_sat [0], r\_sat [1], r\_sat [2], v\_sat [0], v\_sat [1], v\_sat [2]]) \#position$
	and velocity vectors
8	
9	def force $(X, t)$ : #This function gets velocity and acceleration vectors from position and
	velocity vector
LO	$\mathbf{r} = \mathbf{X}[0:3]$
11	$r_n r m = np.linalg.norm(r)$
12	dr = np.zeros((3))
13	dv = np.zeros((3))
14	dr [:] = X[3:6]
15	$dv$ [:] = (-u0 / r_norm ** 3) * r
16	$X_dot = np.array([dr[0], dr[1], dr[2], dv[0], dv[1], dv[2]])$
17	return X_dot
10	

Next, I use the well known RK4 algorithm of finding k1, k2, k3, and k4 to calculate one step forward in  $\mathbf{r}$  by calling the "force" function I created

3

4

5

<sup>2</sup> def RK4\_algorithm(X, t, dt): #This function calculates 1 time-step forward using Runge-Kutta 4 Integration Scheme

k1 = force(X, t) k2 = force(X + dt + lt) / 2 + t

k2 = force(X + dt \* k1 / 2, t + dt / 2)

Then, I created a function named "RK4" that loops the RK4\_algorithm function to calculate all steps from t = 0 to the desired t\_final. The integrator uses a fixed-time-step size that is defined as t\_final divided by dt which is set by the user.

```
def RK4(X, t_final, dt): #This function loops the RK4_algorithm function to integrate
2
         over the whole time of the simulation
            steps = t_{final} / dt
3
            r_{array} = np.zeros((int(steps), 5))
 4
            for i in range((int(steps))):
                Mean\_anomaly = ((2 * np.pi) * (t\_array[i])) / (T)
6
                Ea = Eccentric_Anomaly(Mean_anomaly)
7
                theta = 2 * np.arctan(np.sqrt((1 + e) / (1 - e)) * np.tan(Ea / 2))
8
                r\_star, t\_update, r\_mag = RK4\_algorithm(X, i, dt)
9
                \mathbf{X} = \mathbf{r\_star}
10
                r\_array\left[\,i\,,\ 0\right]\,=\,X[0]
                r_{array}[i, 1] = X[1]
                r_{array}[i, 2] = X[2]
13
                r_{array}[i, 3] = theta
14
                r_{array}[i, 4] = r_{mag}
15
            r\_update = r\_star
16
            return r_array
```

To compare how my RK4 integrator compares to the conic trajectory, I created another function that calculates the same orbit using classical orbital elements and the orbit formula

$$r = \frac{h^2}{\mu} \frac{1}{1 + e\cos(\theta)} \tag{4}$$

where h is specific angular momentum,  $\mu$  is the gravitational parameter of Earth, e is eccentricity, and  $\theta$  is true anomaly.

1

```
def conic_trajectory (T, t_final , dt): #Calculate conic/theoretical orbit using classical
     2
                                                             orbital elements and the orbit formula
                                                                               steps = t_{final} / dt
     3
                                                                                  r_{conic} = np.zeros((int(steps), 5))
       4
                                                                                  for j in range(int(steps)):
     5
                                                                                                              Mean\_anomaly = ((2 * np.pi) * (t\_array[j])) / (T)
     6
                                                                                                              Ea = Eccentric_Anomaly(Mean_anomaly)
       7
                                                                                                             theta = 2 * np.arctan(np.sqrt((1 + e) / (1 - e)) * np.tan(Ea / 2))
r_conic_mag = (h ** 2 / u0) * (1 / (1 + e * np.cos(theta)))
     8
     9
                                                                                                             X = r\_conic\_mag * (np.cos(raan) * np.cos(w + theta) - np.sin(raan) * np.sin(w) + np.sin(
10
                                                                   + theta) * np.cos(i)) # Position X-component
                                                                                                             Y = r\_conic\_mag * (np.sin(raan) * np.cos(w + theta) + np.cos(raan) * np.sin(w + theta) + np.cos(raan) * np.sin(w + theta) + np.cos(raan) * np.sin(w + theta) + np.cos(w + theta) + np.co
                                                                   + theta) * np.cos(i)) \# Position Y-component
```

12	$Z = r_conic_mag * (np.sin(i) * np.sin(w + theta))$	# Position Z-component
13	$r_{conic}[j, 0:5] = X, Y, Z, theta, r_{conic_mag}$	
14	return r_conic	

To ensure that both the RK4 integrator and the conic-trajectory simulator use the same fixed-steps (because RK4 uses time-steps while the conic trajectory simulator uses true-anomaly steps), I use mean and eccentric anomalies to "normalize" the time and true anomalies being used. This ensures that each time-step the RK4 integrator uses corresponds to the true anomaly step that the conic trajectory is using. This allows me to more easily compare the accuracy of the two simulations as I can directly compare the calculations made at every step by RK4 to the calculations made by the same step of the conic-trajectory.

```
def Eccentric_Anomaly(Me): #From HW1
2
           ratio = 1
3
           if Me < np.pi:
5
               E_i = Me + (e / 2.0)
           if Me \geq np.pi:
6
               E_i = Me - (e / 2.0)
7
       # Step 2) Calculate f(E_i) and f'(E_i) and get ratio f/f'. If |ratio| > tolerance,
        calculate new E_{i+1} = E_i - ratio_i and loop until | ratio | meets tolerance
           while np.abs(ratio) > 1e-12:
9
               f1 = E_i - e * np.sin(E_i) - Me
               f_{2} = 1 - e * np.cos(E_i)
               ratio = f1 / f2
12
               E_i = E_i - ratio
13
           Ecc_anomaly = E_i
14
          return Ecc_anomaly
```

Finally, I can run the RK4 integrator and conic-trajectory simulation and compare their results. For the simulation results shown, the initial position and velocity vectors used classical orbital elements a = 26610.2km (which produces an orbital period, T, equal to 12 hrs),  $i = 45^{\circ}$ ,  $\Omega = 45^{\circ}$ ,  $w = 270^{\circ}$ ,  $\theta = 0^{\circ}$ , and  $e = \frac{-J_3 R_E}{2J_2 a} sin(i)$  (the eccentricity value was taken from the "Frozen orbit solution from HW2 discussion on Canvas") where J3 is the perturbation due to Earth's 3rd zonal harmonic and is equal to  $-2.53243x10^{-6}$ , J2 is the perturbation due to Earth's 2nd zonal harmonic and is equal to  $1.08263x10^{-3}$ , and RE is the average radius of the Earth = 6378km. Total propagation time is 5 times the orbital period t = 5 \* T and time-step dt = 1 (and  $\theta$  corresponding to each time step) [See first chunk of code above the definition of the force function].

orbit = RK4(X, t, dt) #Simulate orbit using RK4 integration

 $_{2}$  conic\_orbit = conic\_trajectory(T, t, dt) #Simulate orbit using orbit formula

print(np.max(np.abs(orbit[:, 4] - conic\_orbit [:, 4]))) #Find the largest deviation in position between RK4 simulation and Conic/Orbit formula simulation





(a) View from above the orbit; No perturbations included. Both orbits calculated by RK4 (blue) and Conic-Trajectory (red) are plotted

(b) View from the side the orbit; No perturbations included. Both orbits calculated by RK4 (blue) and Conic-Trajectory (red) are plotted

On the above plots, both the RK4 and Conic-Trajectory orbits are plotted, however they are nearly identical, so they overlap and only the RK4 orbit is seen. You can check this by running the code yourself. The output of the "print(np.max(np.abs(orbit[:, 4] - conic\_orbit[:, 4])))" line takes the directly compares the orbital position calculated by RK4 and Conic-Trajectory by taking the difference. Instead of printing the entire difference array, I only print the maximum difference calculated which will give us an idea as to how accurate the RK4 integrator is to the theoretical conic-trajectory. The code returns the value "0.0010166317733819596" which means the greatest deviation in position between the two simulated orbits is just 1.01 meters.

So, to answer the question: "Does the error ever grow over 500km?" The answer is most likely "no" even if one were to propagate the orbit (without perturbations) for a much longer time.

## 3 Problem 3

Problem 3 is the same as Problem 2, though we now introduce J2 and J3 perturbations to the orbit calculations.

I basically recycle the code from problem 2 making small edits so the functions and variables do not get mixed up with the non-perturbed calculations. The J2 and J3 perturbing accelerations are calculated in the "force\_perturbed" function that plays the same roll that "force" function does in problem 2. I calculate the X, Y, and Z components of the J2 and J3 perturbing accelerations, then take their magnitudes and add them to the acceleration vector felt by the satellite on the orbit:

```
def force_perturbed(X, t): # Calculates acceleration vector, including J2 and J3
    1
                                                                r = X[0:3]
    2
                                                                r_norm = np.linalg.norm(r)
    3
    4
                                                                 \#J2 Perturbing Acceleration
    5
                                                                    aj2_coefficient = (-3.0 / 2.0) * J2 * (u0 / r_norm ** 2) * (Re / r_norm) ** 2
    6
                                                               \begin{array}{l} a_{j}22 = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [0] / r\_norm) \\ a_{j}2y = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [1] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (3 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (3 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (3 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) \\ a_{j}2z = a_{j}2\_coefficient * (1 - 5 * ((r [2] / r\_norm) **2)) * (r [2] / r\_norm) 
    7
    8
    9
                                                                aj2 = np.array([aj2x, aj2y, aj2z])
10
                                                                 \#J3 Perturbing Acceleration (From Schaub H, Junkins, JL. 2009. "Analytical
 12
                                                 Mechanics of Space Systems" pg. 380)
                                                                     aj3_coefficient = (-1.0 / 2.0) * J3 * (u0 / r_norm ** 2) * (Re / r_norm) ** 3
13
                                                                aj3x = aj3\_coefficient * (5 * (7 * (r[2] / r\_norm) ** 3 - 3 * (r[2] / r\_norm)) * (7 * (r[2] / r\_norm))) * (7 * (r[2] / r\_nor
14
                                                 r[0] / r_norm))
                                                                a_{j}3y = a_{j}3\_coefficient * (5 * (7 * (r [2] / r\_norm) ** 3 - 3 * (r [2] / r\_norm)) * (7 * 
                                                 r[1] / r_norm))
                                                                a_{j}3z = a_{j}3_coefficient * (3 * (10 * (r [2] / r_norm) ** 2 - (35.0 / 3.0) * (r [2] /
16
                                                 r_norm) ** 4 - 1))
                                                                aj3 = np.array([aj3x, aj3y, aj3z])
17
18
                                                                dr = np.zeros((3))
19
20
                                                                dv = np.zeros((3))
                                                                dr [:] = X[3:6]
dv [:] = (-u0 / r_n \text{orm} ** 3) * r + aj2 + aj3 \# \text{ACCOUNT FOR J2 AND J3}
21
22
                                                 PERTURBING ACCELERATIONS HERE
23
                                                                X_{dot} = np.array([dr[0], dr[1], dr[2], dv[0], dv[1], dv[2]])
24
                                                                return X_dot
25
```

The same sequence for numerically integrating the orbit with RK4 is followed as in problem 2. Some results of the perturbed RK4 integration are shown below.



(a) Simulation of 10 orbits with J2 and J3 perturbations present, using same COE as problem 2



(c) Simulation of 100 orbits with e = 0.12and the other COE unchanged. J2 and J3 perturbations included.



(b) Simulation of 200 orbits with J2 and J3 perturbations present, using same COE as problem 2



12 (d) Simulation of 100 orbits with e = 0.3J3 and a = 10610.2km keeping the remaining COE unchanged. J2 and J3 perturbations included.

Figures (a) and (b) show the short-term and long-term variation, respectively, in the "frozen orbit" due to J2 and J3 perturbations. The short-term variation in the orbit is minimal and is not distinguishable from the non-perturbed orbit simulation from problem 2<sup>-1</sup>. However, after 200 orbits (though short term variations are negligible) variations in the classical orbital elements begin to manifest as the orbit slightly widens around the equator. Because this orbit has a period of 12 hours, this long-term variation is achieved over  $\approx 100$  days.

Figures (c) and (d) are provided for comparison reasons. Orbit (c) is prop-

<sup>&</sup>lt;sup>1</sup>It should be noted that "short term" is typically defined as on the order of a single orbit. 10 orbits could be considered beyond this range and no longer "short-term," however, if no meaningful variations are seen after 10 orbits, then a time span of less than 10 orbits will also provide no significant variations to the COE

agated for 100 orbits with an eccentricity value slightly larger than (a) and (b). Some variation in the orbit (around the equator) is present, though not dramatic. Orbit (d) has a much smaller semi-major axis at 10610.2 km and an eccentricity value of 0.3. Because this orbit is much closer to the Earth than the others, the effect of J2 and J3 perturbations are much stronger, resulting in major variations in the classical orbital elements after only 100 orbits.

The results of problem 2 and problem 3 follow the basic idea of general perturbation theory. Problem 2 models an unperturbed orbit which can be numerically integrated, but can also be accurately modeled by a algebraic formula - eq. 4. This equation can be used to plot an ellipse (or circle, parabola, and hyperbola in special cases) around a central body. However, it cannot model an orbit subject to perturbing forces. This is where numerical integration comes in handy. By using a fixed-step integration method such as Runge-Kutta 4, Euler Forward, or Newton's method we can include the perturbing forces felt by an orbiting body at each time step and accurately model the orbit with perturbations. One of the downsides of numerical integration is its accuracy and stability dependent on the chosen step-size and can also be computationally intensive and time consuming to simulate.

Some comments: While completing this assignment and thinking about the short-term and long-term variations, I realized that it is a situation of small-things add up over time and can become impactful on the long-term. A common example is if one were to put a few coins or dollars into their piggy bank each day, they wouldn't have much money saved after a few days - but after many years, one could accumulate a (relatively) large chunk of money. So, just because short-term variations are negligible, like for our "frozen" orbit, that does not mean that it will remain "frozen" on the long-term.

Also, on a side note: The RK4 method is much slower than scipy's odeint function. It gets the job done, though.

## 4 Problem 4

In Problem 4, we add atmospheric drag as a perturbation (and remove J2 J3) into the orbit simulation and must find the time it takes a 200km altitude equatorial circular orbit to decay to 50 km<sup>2</sup>.

To add atmospheric drag into our orbit simulation, we must calculate the perturbing force caused by the atmospheric drag on our spacecraft. The magnitude of this perturbing acceleration is dependent on factors such as atmospheric density, the mass, ballistic coefficient, coefficient of drag, and area normal to velocity of the spacecraft.

First, we define the atmospheric density as:

$$\rho = \rho_0 \exp(-\frac{h - h_0}{H}) \tag{5}$$

 $<sup>^2</sup>For$  Problem 4, I used Example\_12\_01.m (located on page e137 here: https://booksite.elsevier.com/9780080977478/downloads/Appendix\_D.pdf) for some guidance

where, from the assignment: "*h* is the instantaneous spacecraft altitude,  $\rho_0$  is a reference density (use  $5.464x10^{-10}km/m^3$ ) at a reference altitude  $h_0$  (use 180 km) and *H* is a scale height (use 29.740 km)."

Next, we define the spacecraft properties: Mass (M) = 250 kg, Area  $(A) = 4m^2$ , and Coefficient of Drag (Cd) = 2.2 which is the standard for low Earth satellites. With these parameters, we can calculate the satellite's Ballistic Coefficient (Bc):

$$B_c = \frac{M}{C_d A} = 28.4091 kg/m^2 \tag{6}$$

The perturbing acceleration due to atmospheric drag is

$$a_{AD} = -\frac{1}{2}\rho V^2 \frac{C_d A}{M} \tag{7}$$

$$= -\frac{1}{2} \frac{\rho \mathbf{V}^2}{B_c} \tag{8}$$

is negative because atmospheric drag is opposes the satellite's velocity vector. V = the relative velocity of the satellite with respect to the Earth's angular velocity vector crossed with the satellite's position vector:

$$\mathbf{V} = \mathbf{V}_{\mathbf{satellite}} - \omega_{\mathbf{Earth}} \times \mathbf{r}_{\mathbf{satellite}} \tag{9}$$

The numerical integration process is nearly identical to the previous methods, with these parameters calculated in the "force\_drag" function. Once the perturbing acceleration is calculated, it is added to eq. 1 which is then propagated through the RK4 algorithm as done previously:

```
def force_drag(X, t): # Calculates acceleration due to atmospheric drag
2
           wE = np.array([0, 0, 7.2921159e-5]) \# Angular Velocity vector of Earth
3
           r = X[0:3]
           r_norm = np.linalg.norm(r)
5
           v = X[3:6]
6
           vrel = v - np.cross(wE, r) # Relative velocity vector of satellite with respect
\overline{7}
        to rotating Earth
           vrel_norm = np.linalg.norm(vrel)
8
           uv = vrel / vrel_norm # Relative velocity unit vectors
9
10
           \# Drag Perturbation
           alt = r_norm - Re \# instantaneous altitude
           # print(alt)
13
           h0 = 180 \# km reference altitude
14
           H = 29.740 \# km scale height
           p = p0 * np.exp(-(alt - h0) / H)
16
17
18
           \# Satellite parameters
           M = 250 \# kg Mass of starlink satellite
19
           A = 4.0 \ \# m^2
20
           Cd = 2.2 \# Drag Coefficient ***This is the standard value for low earth
21
         satellites ***
           Bc = M / (Cd * A) \# Ballistic Coefficient
22
```

20	
24	Fd = -Cd * A / M * p * (1000 * vrel) ** 2 / 2 * uv # Acceleration (yes,
	acceleration not force $> a = f/m$ . Don't mind the variable naming convention) due
	to the atmospheric drag force. 1000 $*$ vrel to get m/s <sup>2</sup>
25	ad = Fd / 1000 $\#$ go back to km/s <sup>2</sup>
26	
27	dr = np.zeros((3))
28	dv = np.zeros((3))
29	dr [:] = X[3:6]
30	$dv [:] = ((-u0 / r_norm ** 3) * r) + ad # ACCOUNT FOR ATMOSPHERIC$
	DRAG ACCELERATION HERE#
31	
32	$X_{dot} = np.array([dr [0], dr [1], dr [2], dv [0], dv [1], dv [2]])$
33	return X_dot

With the parameters defined above, and dt = 1 the code calculates the time to decay down to 50 km  $\approx$  89973 seconds.



Figure 3: Plot displaying the decrease in altitude from a 200km equatorial circular orbit down to 50km due to atmospheric drag

Finally, we investigate how a 10% uncertainty in out reference density value  $\rho_0$  impacts the decay time with a Monte Carlo Analysis.

```
def Monte_Carlo_Decay(p0_original):
1
      uncertainty = 0.1 #10 \%
2
      num_of_runs = 40
3
      decay_times = np.zeros((num_of_runs, 3))
4
      for n in range(num_of_runs):
          p0 = np.random.normal(p0_original, p0_original * uncertainty)
6
          decay_time = Orbit_Decay(p0)
7
          decay\_times[n,\ 0]\ = p0
8
          decay_times[n, 1] = decay_time[-1]
```

```
\begin{array}{ll} decay\_times[n,\ 2]\ =\ n\ +\ 1\\ \#print("p0\ =\ "\ +\ str(decay\_times[n,\ 0])) \end{array}
10
           \#print(decay_times[n, 0], decay_times[n, 1], decay_times[n, 2])
12
13
       #Calculate Uncertainty in Decay Time#
14
       avg_decay_time = np.average(decay_times[:, 1]) #Take average of the decay time
15
16
       deviations = np.zeros((num_of_runs))
       for j in range
(num_of_runs): #Get the deviation of each decay time with respect to
        the average value
           deviations[j] = np.abs(decay_times[j, 1] - avg_decay_time)
18
       uncertainty = np.average(deviations [:]) #Take average of deviations
19
       print("The uncertainty in time to decay with 10'%' uncertainty in p0 is: " + str(
20
        uncertainty))
21
       \#Plot Monte Carlo Results\#
       ax1 = plt.subplot(2, 1, 1)
23
24
       ax1.plot(decay_times[:, 0], decay_times[:, 1])
       plt.xlabel("Time to Decay (sec)")
25
       plt.ylabel("p0")
26
       plt. title ("Monte Carlo Simulation for p0 and Impact on Decay Time")
27
28
       ax2 = plt.subplot(2, 1, 2)
29
       ax2.plot(decay_times[:, 2], decay_times[:, 1])
30
       plt.ylabel("Time to Decay (sec)")
31
       plt.xlabel("Simulations Ran")
       plt. title ("Monte Carlo Decay Time vs. Simulation Runs")
33
34
       plt.grid()
35
       plt.show()
36
37
       return decay_times, uncertainty
38
```

Running the Monte Carlo simulation 40 times with a 10% uncertainty in  $\rho_0$  and dt = 0.5 results in an uncertainty in decay time (down to 50km) of  $\approx 4736$  seconds.



Figure 4: Variation in Decay time for each Monte Carlo simulation ran

Time	to	decay	to	50	km	with	p0	=	(5.110195843565214e-10) is: [91323.0] seconds
Time	to	decay	to	50	km	with	p0	=	(5.829992807952333e-10) is: [84384.0] seconds
Time	to	decay	to	50	km	with	p0	=	(6.494050050939299e-10) is: [75916.0] seconds
Time	to	decay	to	50	km	with	р0	=	(5.065688617189978e-10) is: [91323.0] seconds
Time	to	decay	to	50	km	with	р0	=	(5.432120318002591e-10) is: [90507.5] seconds
Time	to	decay	to	50	km	with	p0	=	(6.066964977891941e-10) is: [81172.0] seconds
Time	to	decay	to	50	km	with	p0	=	(5.534738643705791e-10) is: [88854.0] seconds
Time	to	decay	to	50	km	with	p0	=	(5.204534010987316e-10) is: [91323.0] seconds
Time	to	decay	to	50	km	with	p0	=	(5.847943322780693e-10) is: [84140.0] seconds
Time	to	decay	to	50	km	with	p0	=	(4.756628649092569e-10) is: [91323.0] seconds
Time	to	decay	to	50	km	with	p0	=	(4.889232402054714e-10) is: [91323.0] seconds
Time	to	decay	to	50	km	with	p0	=	(4.1845879304052393e-10) is: [91323.0] seconds
Time	to	decay	to	50	km	with	p0	=	(6.418914417789473e-10) is: [76770.5] seconds
Time	to	decay	to	50	km	with	p0	=	(5.604553051063479e-10) is: [87765.5] seconds
Time	to	decay	to	50	km	with	p0	=	(6.228551597161905e-10) is: [79062.5] seconds
Time	to	decay	to	50	km	with	p0	=	(5.144122116056865e-10) is: [91323.0] seconds
Time	to	decay	to	50	km	with	p0	=	(5.410612176373871e-10) is: [90837.5] seconds
Time	to	decay	to	50	km	with	p0	=	(5.46122573160638e-10) is: [90022.5] seconds
Time	to	decay	to	50	km	with	p0	=	(4.459585852292013e-10) is: [91323.0] seconds
Time	to	decay	to	50	km	with	р0	=	(5.969952779739625e-10) is: [82470.5] seconds
Time	to	decay	to	50	km	with	р0	=	(5.044058518606537e-10) is: [91323.0] seconds
Time	to	decay	to	50	km	with	p0	=	(5.791344018339641e-10) is: [84975.5] seconds
Time	to	decay	to	50	km	with	p0	=	(4.986120038086048e-10) is: [91323.0] seconds
Time	to	decay	to	50	km	with	p0	=	(5.988889320218816e-10) is: [82207.0] seconds
The ι	ince	ertain	ty :	in 1	time	e to d	dec	ay	with 10'%' uncertainty in p0 is: 4735.832499999999
[Fini	ishe	ed in 3	321:	1.3	5]				

Figure 5: Output of Monte Carlo simulation ran 40 times with an uncertainty in the reference atmospheric density  $\rho_0$  of 0.1 and time step = 0.5. Note the run time...

# A Full Code

1	# Pierce Jackson
2	# Applied Orbital Mechanics, Davide Guzzetti, September 30, 2020
3	# HW 3
4	import numpy as np
5	import scipy as sci
6	import scipy.stats as stats
7	import scipy.integrate $\#$ ode solver solve_ivp(function, t_span, y0) : tspan is interval
	of integration
8	from matplotlib import pyplot as plt
9	from mpl_toolkits.mplot3d import Axes3D
10	## Problem 1 $-$ Create algorithms to convert from Classical Orbital Elements $>$
	Position and Velocity Vectors (COE2RV) and vice versa (RV2COE) $##$
11	
12	
13	def COE2RV(a, e, i, RAAN, w, ta):
14	# Used https://web.archive.org/web/20160418175843/https://ccar.colorado.edu/
	asen5070/handouts/cart2kep2002.pdf as a reference
15	
16	r = (a * (1 - e * 2)) / (1 + e * np.cos(ta)) # Get position from orbit formula
17	h = np.sqrt(mu * a * (1 - e * 2)) # Magnitude of specific angular momentum
18	X = r * (np.cos(RAAN) * np.cos(w + ta) - np.sin(RAAN) * np.sin(w + ta) * np.cos(i)
	) # Position X-component
19	Y = r * (np.sin(RAAN) * np.cos(w + ta) + np.cos(RAAN) * np.sin(w + ta) * np.cos(i
	)) # Position Y-component

```
Z = r * (np.sin(i) * np.sin(w + ta)) # Position Z-component
20
            p = a * (1 - e * 2) \# Semilatus Rectum
21
            X_{dot} = ((X * h * e) / (r * p)) * np.sin(ta) - ((h / r) * (np.cos(RAAN) * np.sin(w)))
22
               + ta) + np.sin(RAAN) * np.cos(w + ta) * np.cos(i))) # Velocity X-component
            Y_{dot} = ((Y * h * e) / (r * p)) * np.sin(ta) - ((h / r) * (np.sin(RAAN) * np.sin(w + p)) + (h + e)) + (h + e) + 
23
               + ta) - np.cos(RAAN) * np.cos(w + ta) * np.cos(i))) # Velocity Y-component
            Z_{dot} = ((Z * h * e) / (r * p)) * np.sin(ta) + ((h / r) * (np.sin(i) * np.cos(w + ta)))
              )) # Velocity Z-component
25
            r_vector_from_COE2RV = np.array([X, Y, Z]) # Put X,Y,Z into an array to create a
26
              vector
            v_vector_from_COE2RV = np.array([X_dot, Y_dot, Z_dot]) # Put X_dot, Y_dot, Z_dot
27
              into an array to create a vector
28
            #print("Position Vector (km)= " + str(r_vector_from_COE2RV))
#print("Velocity Vector (km/s) = " + str(v_vector_from_COE2RV))
29
30
31
            return(r_vector_from_COE2RV, v_vector_from_COE2RV)
33
     def RV2COE(r_vector, v_vector):
34
            \# Used "Orbital Mechanics for Engineering Students" 3rd Edition by Howard D.
35
               Curtis; Also used our class notes
36
37
            r = np.linalg.norm(r_vector) \# Distance
            v = np.linalg.norm(v_vector) \# MAgnitude of Velocity or Speed
38
            v_r = np.dot(v_vector, r_vector) / r \# radial velocity
39
            h_vector = np.cross(r_vector, v_vector) # specific angular momentum vector
40
            h = np.linalg.norm(h_vector) \# magnitude of specific angular momentum
41
42
            i = np.arccos(h_vector[2] / h) \# inclination
            n_vector = np.cross([0, 0, 1], h_vector) \# vector pointing to ascending node
43
            n = np.linalg.norm(n_vector) \# magnitude of n
44
45
            if n_vector [1] > 0:
                   RAAN = np.arccos(n_vector[0] / n) \# Right Ascension of the Ascending node
46
47
             if n_vector[1] < 0:
48
                   RAAN = 2 * np.pi - np.arccos(n_vector[0] / n)
             e_vector = (1 / mu) * ((((v ** 2) - (mu / r)) * (r_vector)) - ((r * v_r) * v_vector))
49
                \# eccentricity vector
            e = np.linalg.norm(e_vector) \# eccentricity
            if e_vector [2] > 0:
51
                   w = np.arccos(np.dot(n_vector, e_vector) / (n * e)) # Argument of periapse
             if e_vector [2] < 0:
53
                   w = 2 * np.pi - np.arccos(np.dot(n_vector, e_vector) / (n * e))
54
             if v_r > 0:
                   ta = np.arccos((np.dot(e_vector, r_vector) / (e * r))) # True anomaly
56
             if v_r < 0:
                   ta = 2 * pi - np.arccos((r_vector / r) * (e_vector / e))
58
            Energy = (v ** 2 / 2) - (mu / r)
59
             if e == 1:
60
                   p = h * 2 / mu \# Semilatus Rectum
61
                   return "Orbit is parabolic. Eccentricity = infinity"
62
63
             else :
                   a = -mu / (2 \ast Energy) # Semi-major Axis
64
65
                   p = a * (1 - e * * 2)
66
            print("Semi-major Axis (km) = " + str(a))
67
            print("Eccentricity = " + str(e))
68
            print("Inclination (deg) = " + str(i * (180 / np.pi)))
69
```

```
print("Right Ascension of the Ascending Node (deg) = " + str(RAAN * (180 / np.pi)))
70
71
       print ("Argument of Periapse (deg)= " + str(w * (180 / np.pi)))
       print("True Anomaly (deg) = " + str(ta * (180 / np.pi)))
72
73
       return [a, e, i, RAAN, w, ta]
74
76
   mu = 398600 \# Gravitational Parameter of Earth km<sup>3</sup>/s<sup>2</sup>
77
   r_vector = np.array([-6045.0, -3490.0, 2500.0])
78
   v_{v} = np.array([-3.457, 6.618, 2.533])
79
80
81
   \#COE = RV2COE(r_vector, v_vector)
82
    #COE2RV(COE[0], COE[1], COE[2], COE[3], COE[4], COE[5])
83
84
85
                                                                                -#
86
    #
       ______
87
   ## Problem 2 - Numerically propagate orbit from HW2 ##
88
89
   # Gravitational Parameters
90
   G = 6.67408e - 20 \# Gravitational Constant in km^3/(kg*s^2)
91
Me = 5.9722e+24 \# Mass of Earth in kg
93 Ms = 11110 # Mass of Satellite --> I have chosen the mass of the Hubble Space
        Telescope
   J2 = 1.08263e - 3
94
_{95} J3 = -2.53243e-6
<sup>96</sup> theta_E = np.deg2rad(15.04 / 3600) \# rotation rate of earth in rads/s
97 Re = 6378 \# radius of Earth (km)
   u0 = 398600  # Standard Gravitational Parameter in km<sup>3</sup>/s<sup>2</sup>
98
99
101 a = 22610.2 \# \text{Semi-major Axis}
102 i = np.deg2rad(45) # Inclination in radians (value in function in degrees)
   raan = np.deg2rad(45) \# Right-Ascension of the Ascending Node in radians (value in
103
        function in degrees)
   w = np.deg2rad(270) # Argument of Perigee in radians (value in function in degrees)
104
105
   \#ta = np.deg2rad(0)
   ta = np.deg2rad(0.00) # True Anomaly in radians (value in function in degrees)
106
   e = ((-J3 * Re) / (2 * J2 * a)) * np.sin(i) # Eccentricity of orbit in degrees TAKEN
107
        FROM DISCUSSION
   \# e = 0.3
108
   apogee = a * (1 + e) \# Apogee radius
109
   perigee = a * (1 - e) \# Perigee radius
110
   v_perigee = np.sqrt(2 * ((u0 / perigee) - (u0 / (2 * a)))) # Velcoity of spacecraft at
        perigee in km/s; from conservation of energy equation
   h = perigee * v_perigee # specific angular momentum of spacecraft in km^2/s
112
   T = ((2 * np.pi) / np.sqrt(u0)) * a ** (3.0 / 2.0) \# period in seconds
113
   T_{hrs} = T / 3600.0
114
    #
116
118
   def Orbit_Integrator_RK4_no_Perturbations(a, e, i, raan, w, ta):
119
       \# Calculate State Vectors r & v using the COE2RV I created above
120
       r_sat , v_sat = COE2RV(a, e, i, raan, w, ta) \ \# Position state vector (km) and
        Velocity state vector (km/s) of satellite in the geocentric equitorial frame (km)
```

```
123
        t = 5 * T \# Total time for propagation
        dt = 1 \# time-step
124
        t_{array} = np.linspace(0, t, t / dt + 1) \# make a time array
125
        X = np.array([r\_sat [0], r\_sat [1], r\_sat [2], v\_sat [0], v\_sat [1], v\_sat [2]]) \#
126
         position and velocity vectors
127
        def force (X, t): # This function gets velocity and acceleration vectors from position
128
          and velocity vector
129
            r = X[0:3]
            r_norm = np.linalg.norm(r)
130
131
            dr = np.zeros((3))
            dv = np.zeros((3))
            dr[:] = X[3:6]
133
            dv[:] = (-u0 / r_norm ** 3) * r
            X_{dot} = np.array([dr[0], dr[1], dr[2], dv[0], dv[1], dv[2]])
135
136
            return X_dot
137
        def RK4_algorithm(X, t, dt): # This function calculates 1 time-step forward using
138
         Runge-Kutta 4 Integration Scheme
139
            k1 = force(X, t)
140
            k2 = force(X + dt * k1 / 2, t + dt / 2)
141
            k3 = force(X + dt * k2 / 2, t + dt / 2)
142
            k4 = force(X + dt * k3, t + dt)
143
144
            r_{-update} = X + (dt / 6) * (k1 + 2 * k2 + 2 * k3 + k4)
145
            t_update = t + dt
146
147
            r_mag = np.linalg.norm(r_update)
148
            return(r_update, t_update, r_mag)
149
        def RK4(X, t_final, dt): # This function loops the RK4_algorithm function to
151
         integrate over the whole time of the simulation
            steps = t_final / dt
153
            r_{array} = np.zeros((int(steps), 5))
            for i in range((int(steps))):
154
155
                 Mean\_anomaly = ((2 * np.pi) * (t\_array[i])) / (T)
                Ea = Eccentric_Anomaly(Mean_anomaly)
156
                 theta = 2 * np.arctan(np.sqrt((1 + e) / (1 - e)) * np.tan(Ea / 2))
                 r\_star\,,\ t\_update,\ r\_mag = RK4\_algorithm(X,\,i,\,dt)
158
                 X = r_star
159
                 r\_array\left[\,i\,,\ 0\right]\,=\,X[0]
160
                 r\_array\,[\,i\,,\ 1]\ = X[1]
161
                 r_{array}[i, 2] = X[2]
162
                 r\_array[i, 3] = theta
                r_{array}[i, 4] = r_{mag}
164
165
            r_update = r_star
            return r_array
166
167
        def Eccentric_Anomaly(Me): # From HW1
168
            ratio = 1
169
170
            if Me < np.pi:
                E_i = Me + (e / 2.0)
172
             if Me \geq np.pi:
                E_i = Me - (e / 2.0)
173
        # Step 2) Calculate f(E_i) and f'(E_i) and get ratio f/f'. If |ratio| > tolerance,
174
```

```
calculate new E_i + 1 = E_i - ratio_i and loop until | ratio | meets tolerance
            while np.abs(ratio) > 1e-12:
                f1 = E_i - e * np.sin(E_i) - Me
                f2 = 1 - e * np.cos(E_i)
177
                ratio\,=\,{\rm f1} / {\rm f2}
178
                E_{-i} = E_{-i} - ratio
179
180
            Ecc_anomaly = E_i
            return Ecc_anomaly
181
182
        def conic_trajectory (T, t_final, dt): # Calculate conic/theoretical orbit using
183
          classical orbital elements and the orbit formula
            steps = t_{final} / dt
184
            r_{\text{conic}} = \text{np.zeros}((\text{int}(\text{steps}), 5))
185
            for j in range(int(steps)):
186
                Mean\_anomaly = ((2 * np.pi) * (t\_array[j])) / (T)
187
                Ea = Eccentric_Anomaly(Mean_anomaly)
188
189
                theta = 2 * np.arctan(np.sqrt((1 + e) / (1 - e)) * np.tan(Ea / 2))
                r_{conic_mag} = (h ** 2 / u0) * (1 / (1 + e * np.cos(theta)))
190
                X = r\_conic\_mag * (np.cos(raan) * np.cos(w + theta) - np.sin(raan) * np.sin(w)
191
          + theta) * np.cos(i)) # Position X-component
                Y = r_conic_mag * (np.sin(raan) * np.cos(w + theta) + np.cos(raan) * np.sin(w)
          + theta) * np.cos(i)) # Position Y-component
                Z = r_{conic_mag} * (np.sin(i) * np.sin(w + theta)) # Position Z-component
194
                r_{conic}[j, 0:5] = X, Y, Z, theta, r_{conic_mag}
            return r_conic
196
        orbit = RK4(X, t, dt) # Simulate orbit using RK4 integration
        conic_orbit = conic_trajectory(T, t, dt) # Simulate orbit using orbit formula
198
199
        #-----#
        u, v = np.mgrid[0: 2 * np.pi: 100j, 0: np.pi: 50j]
201
202
        x\_sphere = Re * np.cos(u) * np.sin(v)
        y\_sphere = Re * np.sin(u) * np.sin(v)
203
        z\_sphere = Re * np.cos(v)
204
205
        fig = plt.figure()
206
        ax = fig.add_subplot(111, projection='3d')
207
208
        ax.set_aspect("equal")
        ax.plot(conic_orbit [:, 0], conic_orbit [:, 1], conic_orbit [:, 2], color="red")
209
        ax.plot(orbit [:, 0], orbit [:, 1], orbit [:, 2], color="darkblue")
210
        ax.plot_surface(x_sphere, y_sphere, z_sphere, rstride=3, cstride=3, color='none',
211
         edgecolor='k', shade=0
        set_axes_equal(ax)
212
        plt.show()
213
214
        \#print((orbit [:30, 4]), (conic_orbit [:30, 4]))
215
        print(np.max(np.abs(orbit[:, 4] - conic_orbit [:, 4]))) # Find the largest deviation
216
         in position between RK4 simulation and Conic/Orbit formula simulation
217
218
    def Orbit_Integrator_RK4_J2_J3(a, e, i, raan, w, ta):
219
220
221
        \# Calculate State Vectors r & v using the COE2RV I created above
        r_sat, v_sat = COE2RV(a, e, i, raan, w, ta) # Position state vector (km) and
222
         Velocity state vector (km/s) of satellite in the geocentric equitorial frame (km)
        t = 100 * T
224
```

```
dt = 10
225
                 X = np.array([r_sat [0], r_sat [1], r_sat [2], v_sat [0], v_sat [1], v_sat [2]])
227
                 def force_perturbed(X, t): # Calculates acceleration vector, including J2 and J3
228
                          r = X[0:3]
229
                         r_{norm} = np.linalg.norm(r)
                          \#J2 Perturbing Acceleration
232
                          aj2_coefficient = (-3.0 / 2.0) * J2 * (u0 / r_norm ** 2) * (Re / r_norm) ** 2
233
                          \begin{array}{l} aj2x = aj2\_coefficient * (1 - 5 * ((r [2] / r\_norm) ** 2)) * (r [0] / r\_norm) \\ aj2y = aj2\_coefficient * (1 - 5 * ((r [2] / r\_norm) ** 2)) * (r [1] / r\_norm) \\ aj2z = aj2\_coefficient * (3 - 5 * ((r [2] / r\_norm) ** 2)) * (r [2] / r\_norm) \\ \end{array} 
234
235
                         aj2 = np.array([aj2x, aj2y, aj2z])
237
238
                          # J3 Perturbing Acceleration (From Schaub H, Junkins, JL. 2009. "Analytical
239
                   Mechanics of Space Systems" pg. 380)
aj3_coefficient = (-1.0 / 2.0) * J3 * (u0 / r_norm ** 2) * (Re / r_norm) ** 3
240
                         aj3x = aj3_coefficient * (5 * (7 * (r [2] / r_norm) ** 3 - 3 * (r [2] / r_norm)) *
241
                    r[0] / r_norm))
                         a_{j}3y = a_{j}3\_coefficient * (5 * (7 * (r [2] / r\_norm) ** 3 - 3 * (r [2] / r\_norm)) * (7 * 
                    r[1] / r_norm))
                         a_{j}3z = a_{j}3_coefficient * (3 * (10 * (r [2] / r_norm) ** 2 - (35.0 / 3.0) * (r [2] /
243
                    r_{norm} ** 4 - 1))
                         aj3 = np.array([aj3x,\ aj3y,\ aj3z])
244
245
                         dr = np.zeros((3))
246
                         dv = np.zeros((3))
247
                         dr[:] = X[3:6]
248
                         dv[:] = (-u0 / r_norm ** 3) * r + aj2 + aj3 \# ACCOUNT FOR J2 AND J3
249
                    PERTURBING ACCELERATIONS HERE
                         X_{dot} = np.array([dr[0], dr[1], dr[2], dv[0], dv[1], dv[2]])
251
                          return X_dot
252
253
                 def RK4_algorithm_perturbed(X, t, dt): # This function calculates 1 time-step
254
                    forward using Runge-Kutta 4 Integration Scheme
255
256
                          k1 = force_perturbed(X, t)
                         k2 = force_perturbed(X + dt * k1 / 2, t + dt / 2)
257
                          k3 = force\_perturbed(X + dt * k2 / 2, t + dt / 2)
258
259
                         k4 = force_perturbed(X + dt * k3, t + dt)
260
                         r_update = X + (dt / 6) * (k1 + 2 * k2 + 2 * k3 + k4)
261
                         t_{-update} = t + dt
262
                          return(r_update, t_update)
263
264
                 def RK4_perturbed(X, t_final, dt): # This function loops the RK4_algorithm function
265
                    to integrate over the whole time of the simulation
                         steps = t_{final} / dt
266
                          r_{array} = np.zeros((int(steps), 4))
267
                          t = 0
268
                          for i in range(int(steps)):
269
270
                                  r_star, t_update, = RK4_algorithm_perturbed(X, t, dt)
                                  \mathbf{X} = \mathbf{r\_star}
271
272
                                  t = t_update
                                  r\_array\left[\,i\,,\ 0\right]\,=\,X[0]
                                  r_{array}[i, 1] = X[1]
274
```

```
r_{array}[i, 2] = X[2]
275
                r_{array}[i, 3] = t
           r_update = r_star
277
           return r_array
278
279
        orbit_perturbed = RK4_perturbed(X, t, dt) \# Simulate perturbed orbit
280
281
        #-----#
282
        u, v = np.mgrid[0: 2 * np.pi: 100j, 0: np.pi: 50j]
283
        x\_sphere = Re * np.cos(u) * np.sin(v)
284
        y\_sphere = Re * np.sin(u) * np.sin(v)
285
286
        z\_sphere = Re * np.cos(v)
287
        fig = plt.figure()
288
        ax = fig.add_subplot(111, projection='3d')
289
        ax.set_aspect("equal")
290
        ax.plot(orbit_perturbed [:, 0], orbit_perturbed [:, 1], orbit_perturbed [:, 2], color="
         darkblue")
        ax.plot_surface(x_sphere, y_sphere, z_sphere, rstride=3, cstride=3, color='none',
292
         edgecolor='k', shade=0)
        set_axes_equal(ax)
293
        plt.show()
294
295
296
   \#p0 = 5.464e - 10 \ \# \text{ km/m}^3 reference density; Must be outside of function for monte
297
         carlo to work
298
299
   def Orbit_Decay(p0):
300
        a_d = Re + 200
301
        e_d = np.deg2rad(0)
302
303
        i_d = np.deg2rad(0)
        raan_d = np.deg2rad(50)
304
305
        w_d = np.deg2rad(90)
306
        ta_d = np.deg2rad(10)
307
        # Calculate State Vectors r & v using the COE2RV I created above
        r_sat, v_sat = COE2RV(a_d, e_d, i_d, raan_d, w_d, ta_d) # Position state vector (km)
308
         and Velocity state vector (km/s) of satellite in the geocentric equitorial frame (
         km)
        T = ((2 * np.pi) / np.sqrt(mu)) * a_d ** (3.0 / 2.0)
309
        t = 17.2 * T
310
        dt = 0.5
311
        X = np.array([r_sat [0], r_sat [1], r_sat [2], v_sat [0], v_sat [1], v_sat [2]])
312
313
        def force_drag(X, t): # Calculates acceleration due to atmospheric drag
314
315
            wE = np.array([0, 0, 7.2921159e-5]) \# Angular Velocity vector of Earth
316
317
           r = X[0:3]
           r_norm = np.linalg.norm(r)
318
            v = X[3:6]
319
            vrel = v - np.cross(wE, r) # Relative velocity vector of satellite with respect
         to rotating Earth
321
            vrel_norm = np.linalg.norm(vrel)
            uv = vrel / vrel_norm # Relative velocity unit vectors
            # Drag Perturbation
            alt = r_norm - Re \# instantaneous altitude
```

```
# print(alt)
            h0 = 180 \# km reference altitude
327
            H = 29.740 \ \# \text{ km} \text{ scale height}
328
            p = p0 * np.exp(-(alt - h0) / H)
329
            # Satellite parameters
332
            M = 250 \# kg Mass of starlink satellite
            A = 4.0 \ \# m^2
            Cd = 2.2 \# Drag Coefficient ***This is the standard value for low earth
334
          satellites ***
            Bc = M / (Cd * A) \# Ballistic Coefficient
            Fd = -Cd * A / M * p * (1000 * vrel) ** 2 / 2 * uv # Acceleration (yes,
337
         acceleration not force --> a = f/m. Don't mind the variable naming convention) due
          to the atmospheric drag force. 1000 \ast vrel to get m/s^2
            ad = Fd / 1000 \# go back to km/s<sup>2</sup>
338
            dr = np.zeros((3))
340
341
            dv = np.zeros((3))
            dr [:] = X[3:6]
342
            dv [:] = ((-u0 / r_norm ** 3) * r) + ad # ACCOUNT FOR ATMOSPHERIC
343
         DRAG ACCELERATION HERE#
345
            X_{dot} = np.array([dr[0], dr[1], dr[2], dv[0], dv[1], dv[2]])
            return X_dot
346
347
        def RK4_algorithm_drag(X, t, dt): # This function calculates 1 time-step forward
348
         using Runge-Kutta 4 Integration Scheme
349
            k1 = force_drag(X, t)
            k2 = force_drag(X + dt * k1 / 2, t + dt / 2)
351
            k3 = force_drag(X + dt * k2 / 2, t + dt / 2)
352
            k4 = force\_drag(X + dt * k3, t + dt)
353
354
            r_update = X + (dt / 6) * (k1 + 2 * k2 + 2 * k3 + k4)
355
356
            t_update = t + dt
            return(r_update, t_update)
357
358
        def RK4_drag(X, t_final, dt): # This function loops the RK4_algorithm function to
359
         integrate over the whole time of the simulation
            steps = t_final / dt
360
            r_{array} = np.zeros((int(steps), 5))
361
362
            t\ = 0
            for i in range(int(steps)):
363
                r_star, t_update, = RK4_algorithm_drag(X, t, dt)
364
365
                \mathbf{X} = \mathbf{r\_star}
366
                t = t_update
367
                r\_array\left[\,i\,,\ 0\right]\,=\,X[0]
368
                r_{array}[i, 1] = X[1]
369
                r_{array}[i, 2] = X[2]
                r\_array\,[\,i\,,\ 3]\,=t
                r_{array}[i, 4] = np.sqrt(X[0] ** 2 + X[1] ** 2 + X[2] ** 2) - Re
                if 49.95 \leq r_{array}[i, 4] \leq 50.05: # We want to know the time when
         altitude = 50km, so stop loop once altitude = 50km
                    return r_array
374
                    break
```

```
#print(r_array[i, 4])
                       r_update = r_star
377
                       return r_arrav
378
379
                orbit_decay = RK4_drag(X, t, dt) # Simulate perturbed orbit
380
381
                decay_time = orbit_decay[:, 3]
                decay_time = np.trim_zeros(decay_time)
383
384
385
                #plt.rcParams.update({'font.size': 15})
                #plt.plot(decay_time, orbit_decay [: len(decay_time), 4])
386
                #plt.xlabel('Time (sec)')
387
                #plt.ylabel('Altitude [m]')
388
                #plt. title ('Time Evolution of Altitude with Atmospheric Drag')
389
                #plt.grid()
390
                #plt.show()
391
392
                print("Time to decay to 50 km with p0 = (" + str(p0) + ") is: [" + str(decay_time
393
                  [-1] + "] seconds")
394
                return decay_time
395
                #-----#
396
                u, v = np.mgrid[0: 2 * np.pi: 100j, 0: np.pi: 50j]
397
398
                x\_sphere = Re * np.cos(u) * np.sin(v)
                y\_sphere = Re * np.sin(u) * np.sin(v)
399
                z\_sphere = Re * np.cos(v)
400
401
                fig = plt.figure()
402
                ax = fig.add_subplot(111, projection='3d')
403
                ax.set_aspect("equal")
404
                ax.plot(orbit_decay [:, 0], orbit_decay [:, 1], orbit_decay [:, 2], color="darkblue")
405
                ax. plot\_surface (x\_sphere, y\_sphere, z\_sphere, rstride=3, cstride=3, color='none', as a sphere and a spher
406
                  edgecolor='k', shade=0)
407
                set_axes_equal(ax)
                # plt.show()
408
409
410
411
       def Monte_Carlo_Decay(p0_original):
                uncertainty = 0.1 \# 0.6827 \# 1 sigma variation
412
                num_of_runs = 1
413
                decay\_times = np.zeros((num\_of\_runs, 3))
414
                for n in range(num_of_runs):
415
                       p0 = np.random.normal(p0_original, p0_original * uncertainty)
416
                       decay_time = Orbit_Decay(p0)
417
                       decay_times[n, 0] = p0
418
                       decay_times[n, 1] = decay_time[-1]
419
                       decay_times[n, 2] = n + 1
420
                       \#print("p0 = " + str(decay_times[n, 0]))
421
                       #print(decay_times[n, 0], decay_times[n, 1], decay_times[n, 2])
422
423
                #Calculate Uncertainty in Decay Time#
424
                avg_decay_time = np.average(decay_times[:, 1]) #Take average of the decay time
425
426
                deviations = np.zeros((num_of_runs))
                for j in range(num_of_runs): #Get the deviation of each decay time with respect to
427
                  the average value
                       deviations[j] = np.abs(decay_times[j, 1] - avg_decay_time)
428
                uncertainty = np.average(deviations [:]) #Take average of deviations
429
```

```
print ("The uncertainty in time to decay with 10'%' uncertainty in p0 is: " + str(
430
         uncertainty))
431
        \#Plot Monte Carlo Results\#
432
        ax1 = plt.subplot(2, 1, 1)
433
        ax1.plot(decay_times[:, 0], decay_times[:, 1])
434
        plt.xlabel("Time to Decay (sec)")
435
        plt.ylabel("p0")
436
        plt. title ("Monte Carlo Simulation for p0 and Impact on Decay Time")
437
438
        ax2 = plt.subplot(2, 1, 2)
439
        ax2.plot(decay_times [:, 2], decay_times [:, 1])
440
        plt.ylabel("Time to Decay (sec)")
441
        plt.xlabel("Simulations Ran")
442
        plt. title ("Monte Carlo Decay Time vs. Simulation Runs")
443
444
445
        plt.grid()
        plt.show()
446
447
        return decay_times, uncertainty
448
449
       -----#
    #
450
451
452
    def set_axes_equal(ax):
453
         " Make axes of 3D plot have equal scale so that spheres appear as spheres,
454
        cubes as cubes, etc.. This is one possible solution to Matplotlib's
455
        ax.set_aspect ('equal') and ax.axis ('equal') not working for 3D.
456
457
        Input
458
        ax: a matplotlib axis, e.g., as output from plt.gca().
459
460
461
        x_{limits} = ax.get_{xlim3d}()
462
        y_{limits} = ax.get_{ylim3d}()
463
464
        z_{\text{limits}} = ax.get_{z}lim3d()
465
466
        x_range = abs(x_limits[1] - x_limits[0])
        x_{middle} = np.mean(x_{limits})
467
        y_range = abs(y_limits[1] - y_limits[0])
468
        y_middle = np.mean(y_limits)
469
        z_{\text{range}} = abs(z_{\text{limits}}[1] - z_{\text{limits}}[0])
470
471
        z_middle = np.mean(z_limits)
472
        \# The plot bounding box is a sphere in the sense of the infinity
473
        \# norm, hence I call half the max range the plot radius.
474
        plot_radius = 0.5 * max([x_range, y_range, z_range])
475
476
        ax.set_xlim3d([x_middle - plot_radius, x_middle + plot_radius])
477
        ax.set_ylim3d([y_middle - plot_radius, y_middle + plot_radius])
478
        ax.set_zlim3d([z_middle - plot_radius, z_middle + plot_radius])
479
480
481
    #Orbit_Integrator_RK4_no_Perturbations(a, e, i, raan, w, ta)
482
483
    #Orbit_Integrator_RK4_J2_J3(a, e, i, raan, w, ta)
    #Orbit_Decay(5.464e-10)
484
485 Monte_Carlo_Decay(5.464e-10)
```